

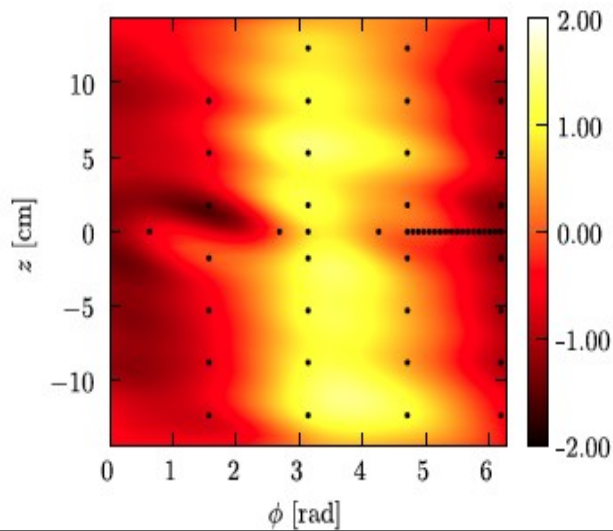
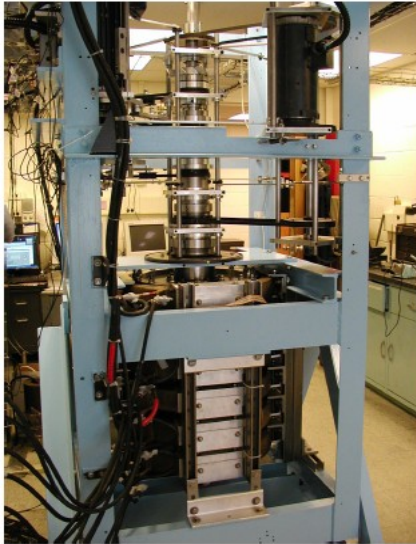
Magnetized spherical Couette flow

Christophe Gissinger, Eric Edlund, Austin Roach, Erik
Spence, Hantao Ji, Jeremy Goodman

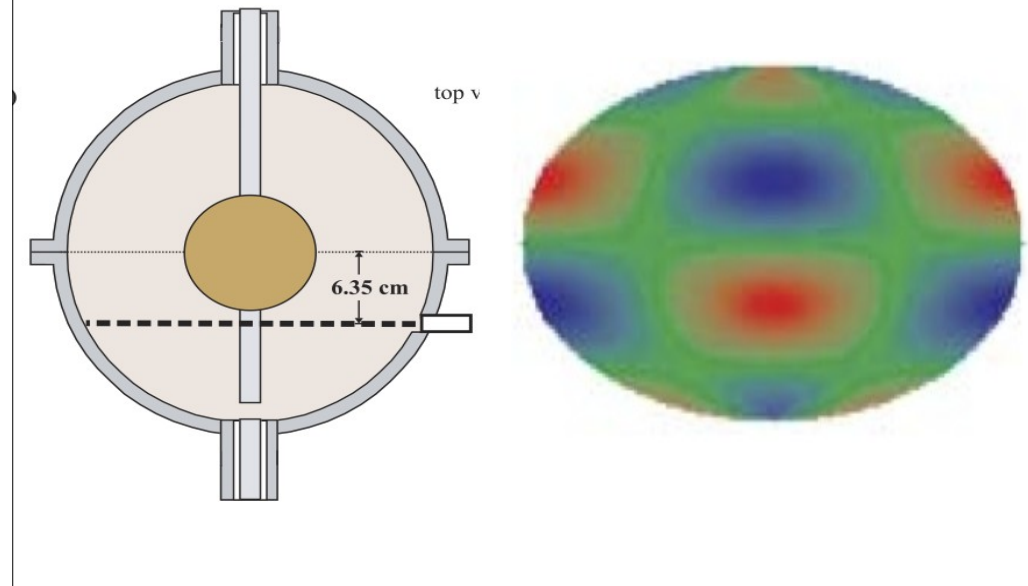
Princeton University, NJ, USA

Context

Princeton MRI experiment



Maryland experiment

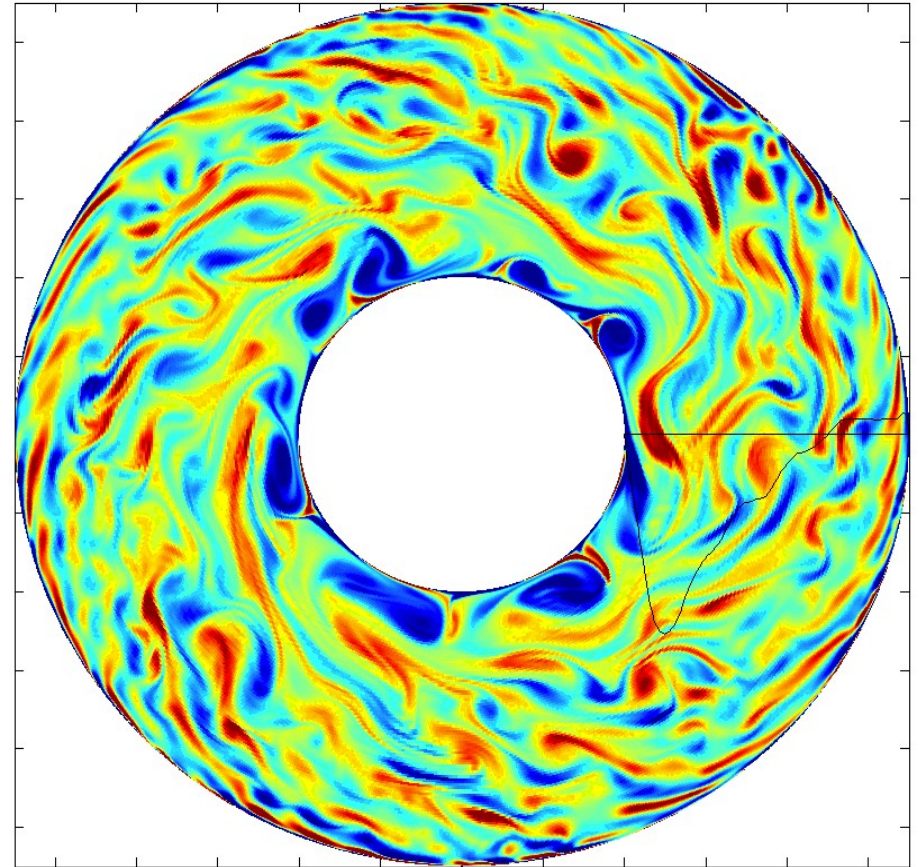


- Non-axisymmetric MHD instabilities
- $m=1$ or $m=2$ azimuthal wavenumber
- Symmetric and antisymmetric modes

Numerical Code

PaRoDy (Dormy et al) :

- Spherical geometry
- *Poloidal-Toroidal* decomposition
- Finite differences in radius
- Cranck-Nicholson / Adams Bashford time-stepping .



Insulating / conducting / ferromagnetic spheres

Full sphere or Couette problem

No slip / Stress Free BC for velocity

Full MHD problem

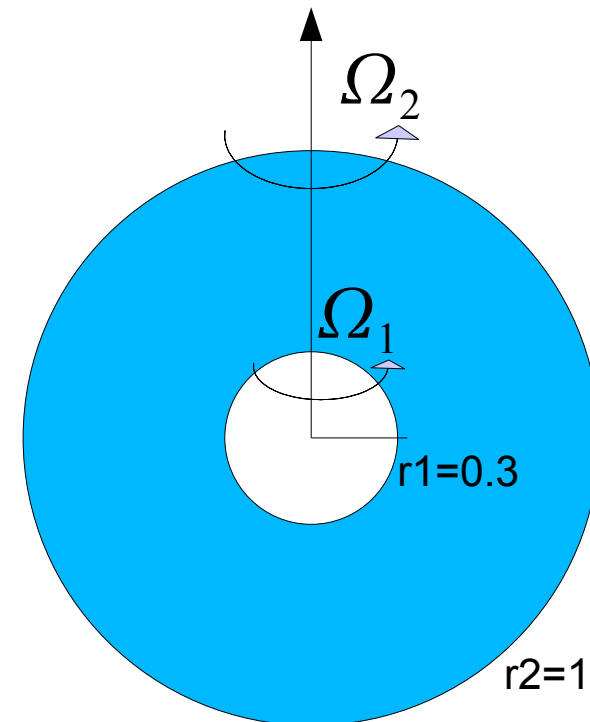
Navier-Stokes equation:

$$\rho \frac{\partial \mathbf{u}}{\partial t} + \rho (\mathbf{u} \nabla) \mathbf{u} = -\nabla P + \rho \nu \nabla^2 \mathbf{u} + \mathbf{F} + \mathbf{j} \times \mathbf{B}$$

Induction equation:

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{u} \times \mathbf{B}) + \frac{1}{\mu_0 \sigma} \nabla^2 \mathbf{B}$$

- Insulating outer and inner spheres
- Dipolar or vertical magnetic field imposed
- With or without overall rotation :
 - $\Omega_1 = 8 \Omega_2$
 - $\Omega_2 = 0$

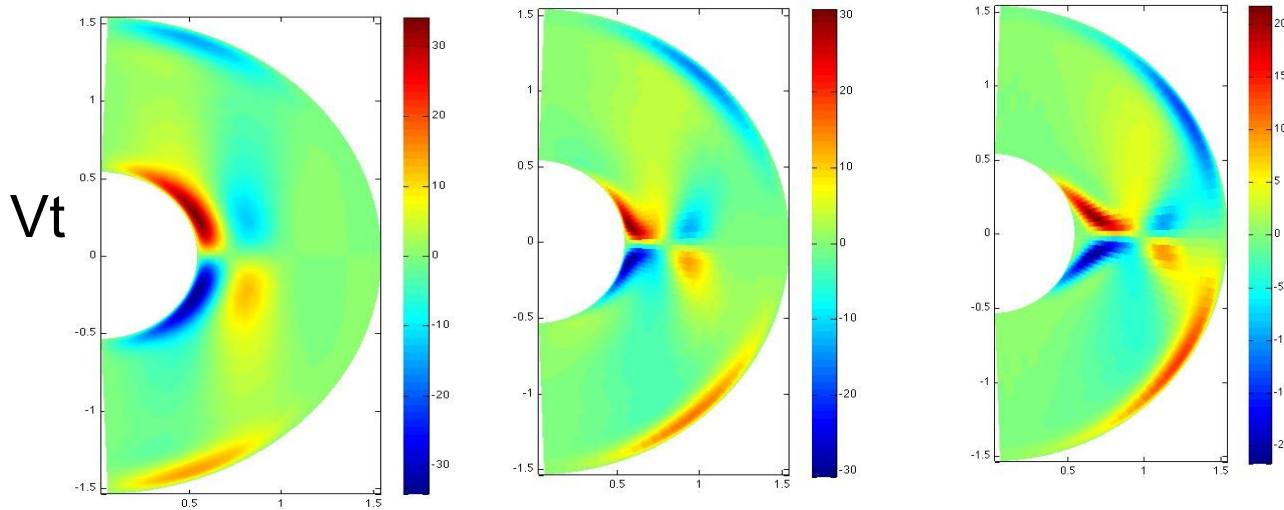
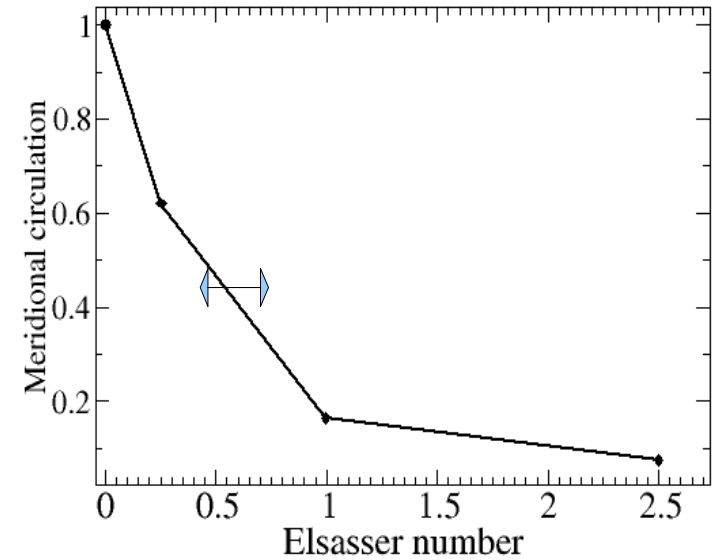
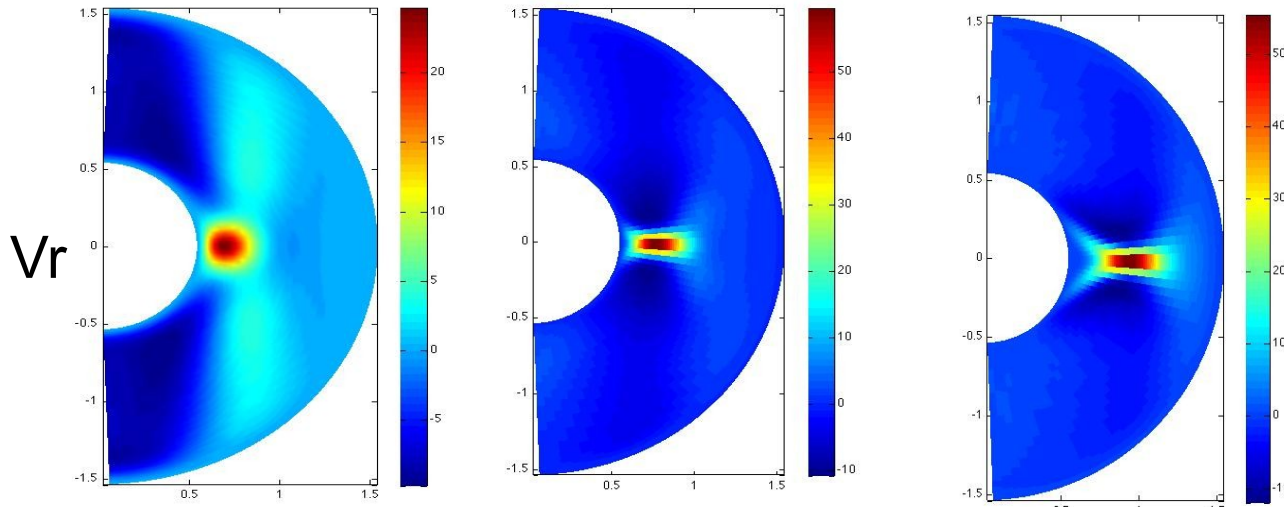


$$\text{Re} = \frac{(\Omega (r_2 - r_1)^2)}{\nu}$$

$$\Lambda = \frac{B_0}{\sqrt{(\rho \mu \eta \Omega)}}$$

Suppression of meridional circulation

Re=330



=0

=1

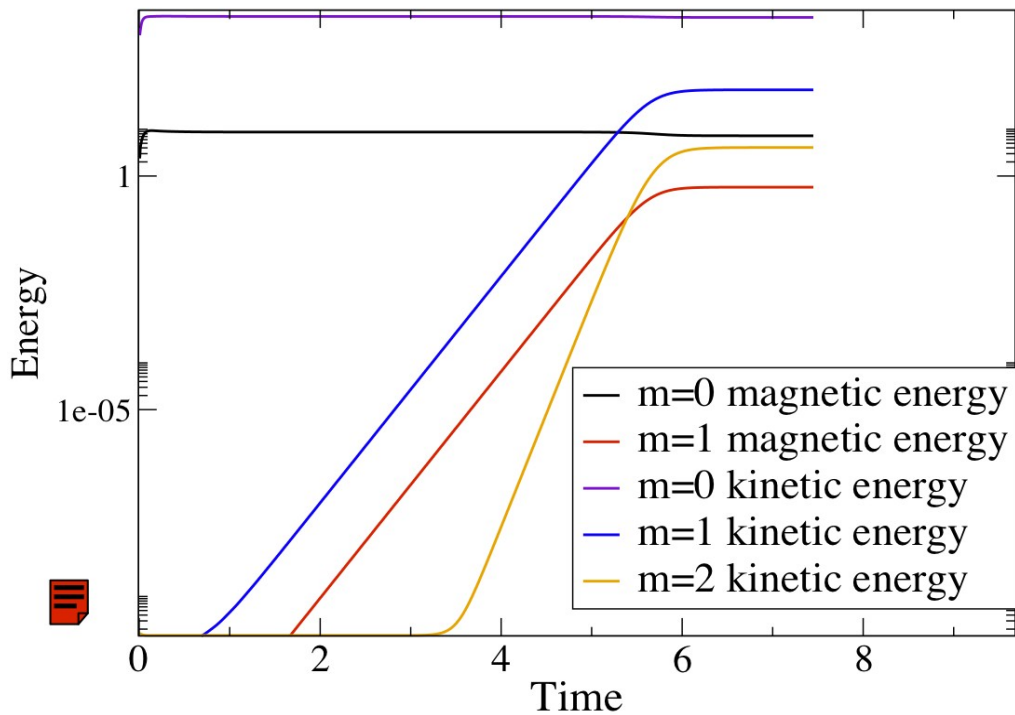
=2.5

Inhibition of the meridional circulation by magnetic field

Non-axisymmetric MHD instability

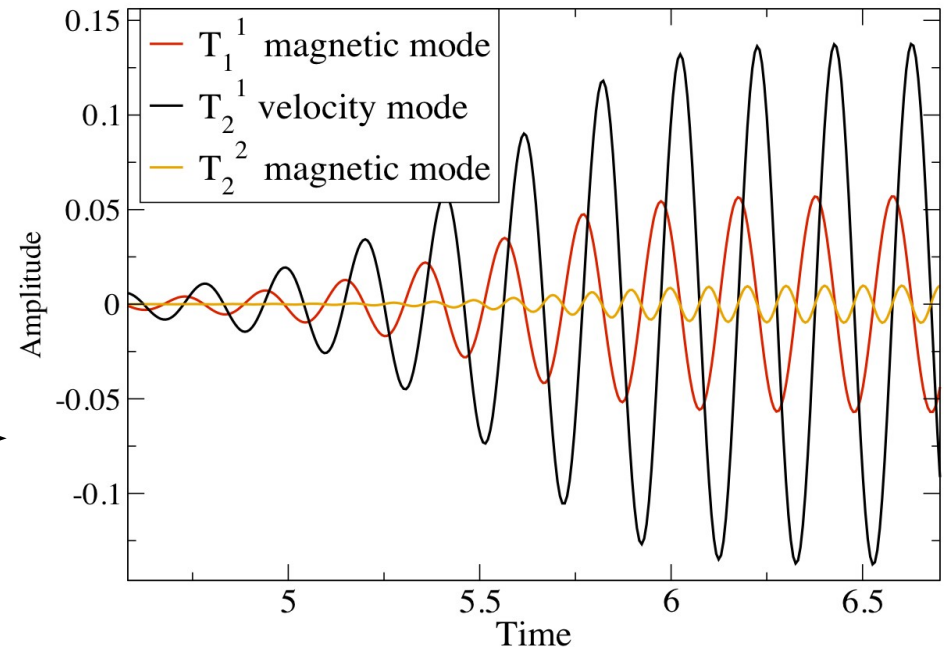
Re=330

$\beta=0.5$



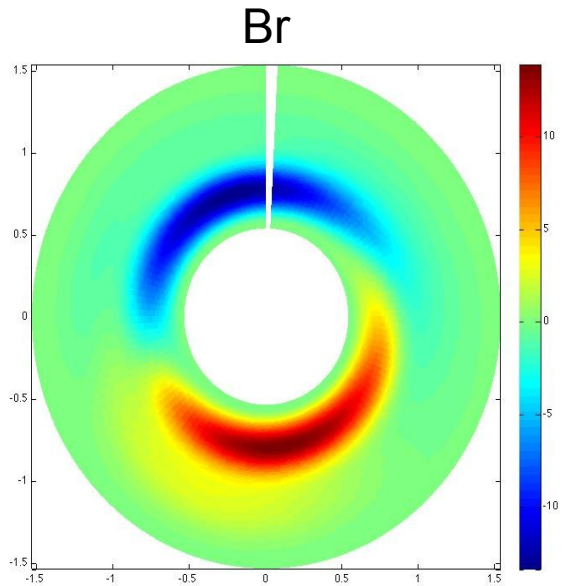
→ m=1 dominated instability
→ smaller m=2 component

Oscillating structure

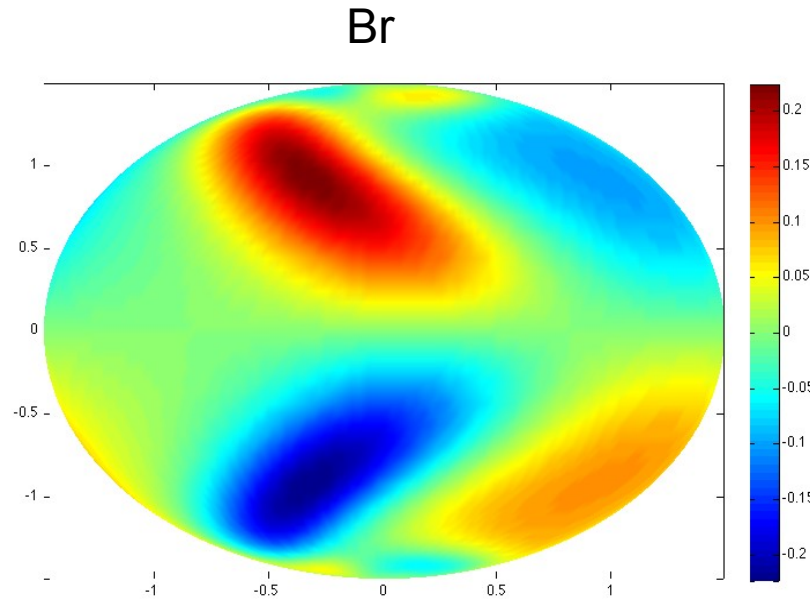


Structure of the instability

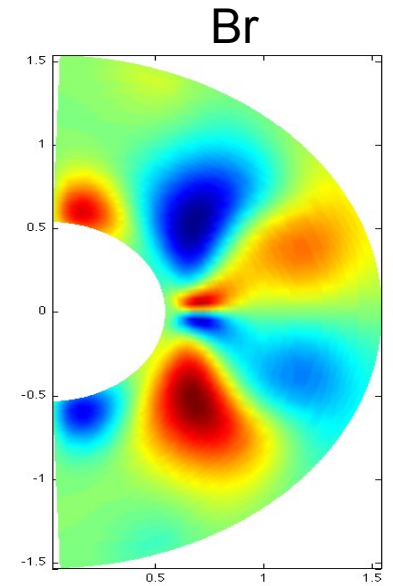
$Re=330$ $\nu=0.5$



Equatorial plane



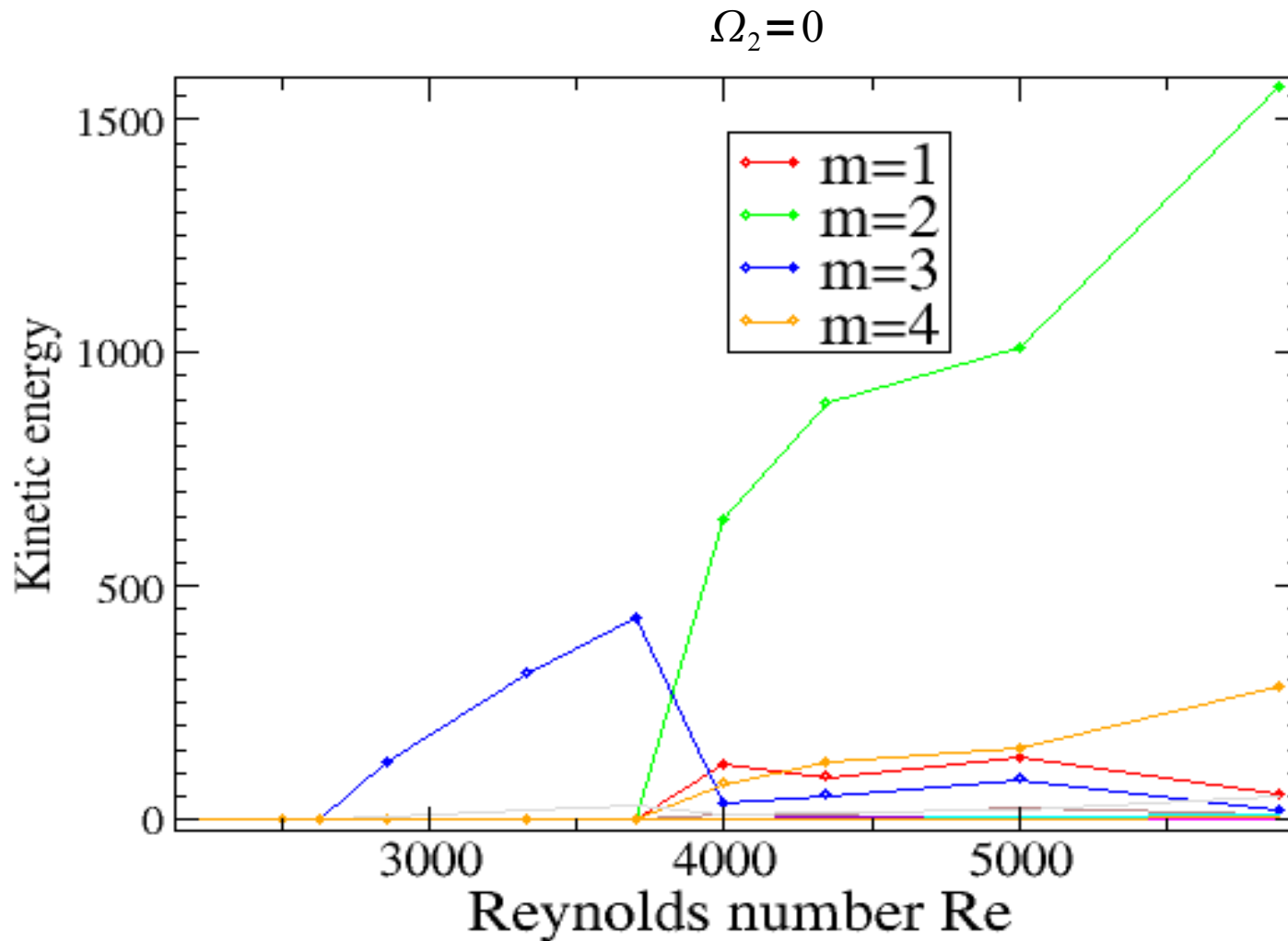
Close to the outer sphere



(R,Z) plane

- Antisymmetric $m=1$ magnetic field
- Instability occurs for dipolar and vertical field
- Instability occurs with or without overall rotation

Hydrodynamical instabilities at Large Re

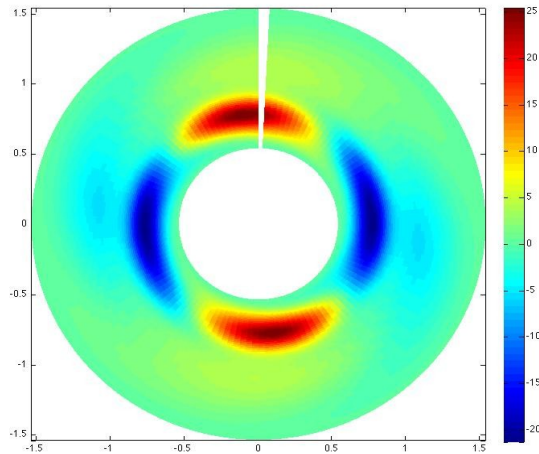


Without rotation: Kelvin-Helmoltz destabilization of the equatorial jet

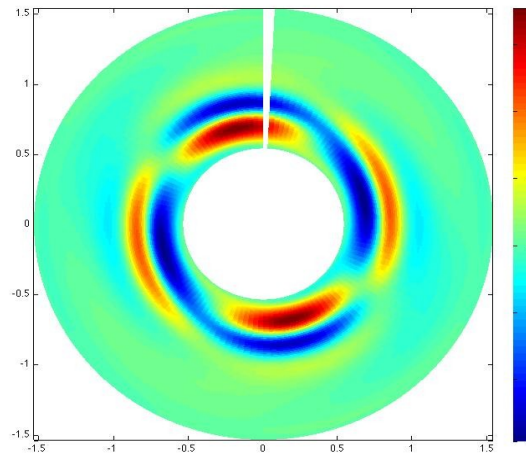
Strong rotation: destabilization of the Stewartson layer

hydrodynamical instability

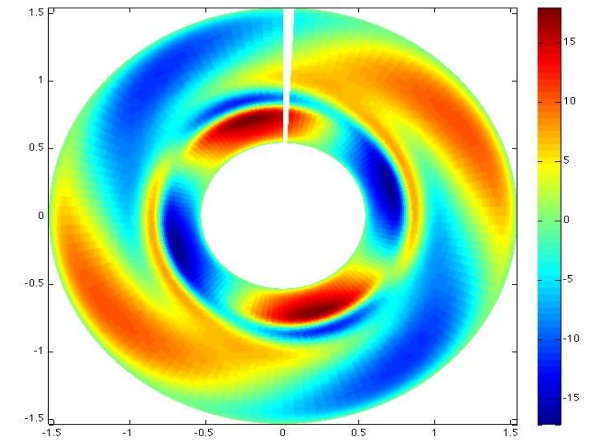
$$\Omega_2 = \cdot$$



V_r

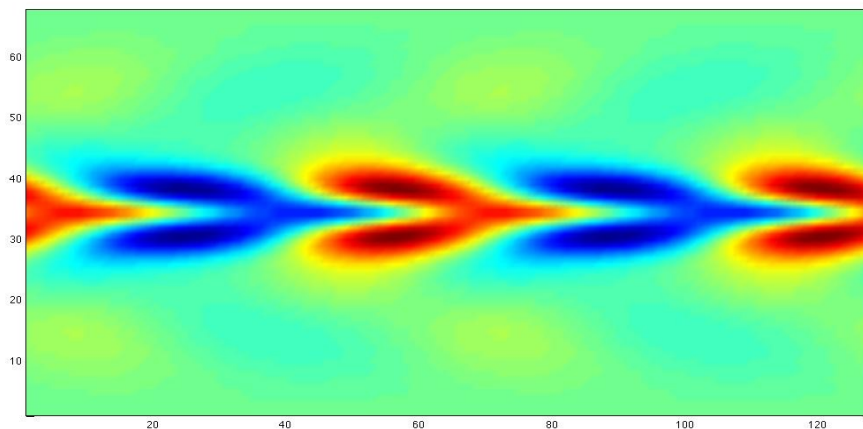


V_t



V_p

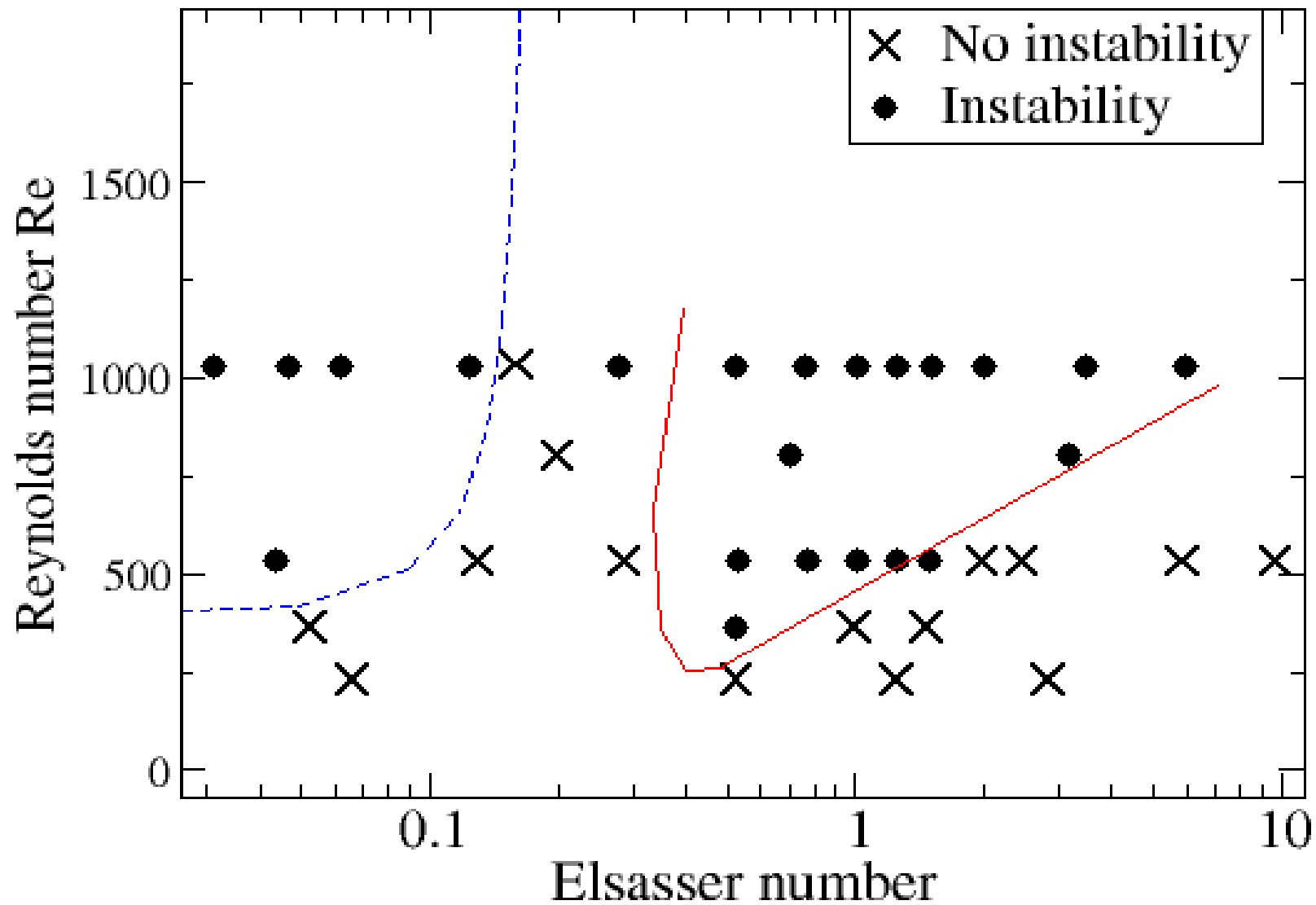
V_r in (Z, ϕ) plane :



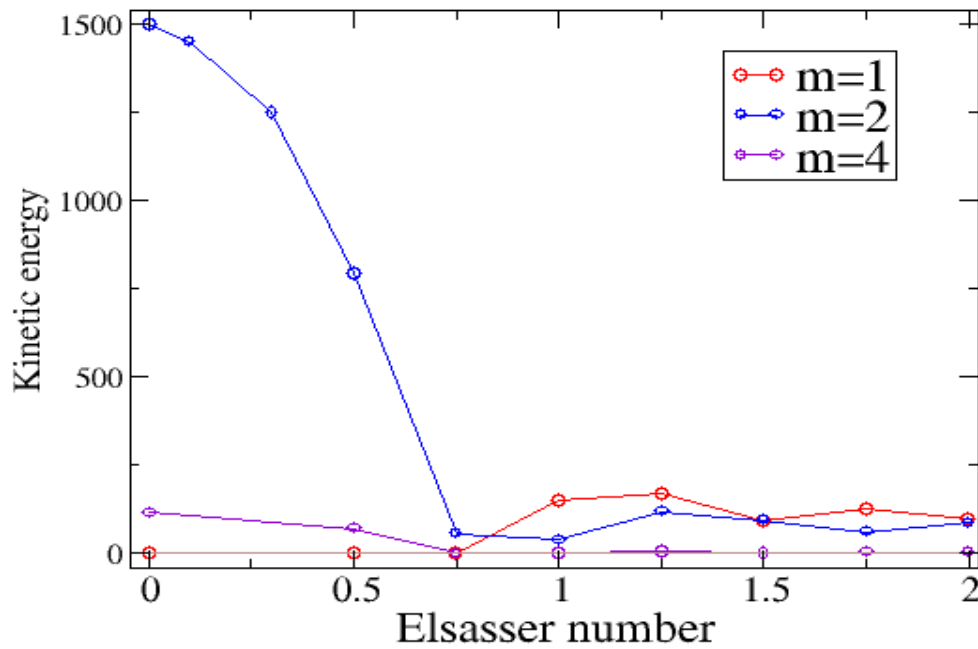
- ▶ $m=2$ non-axisymmetric instability
- ▶ Kelvin-Helmoltz instability of equatorial jet
- ▶ Equatorially symmetric structure
- ▶ Suppressed by magnetic field

Parameter space

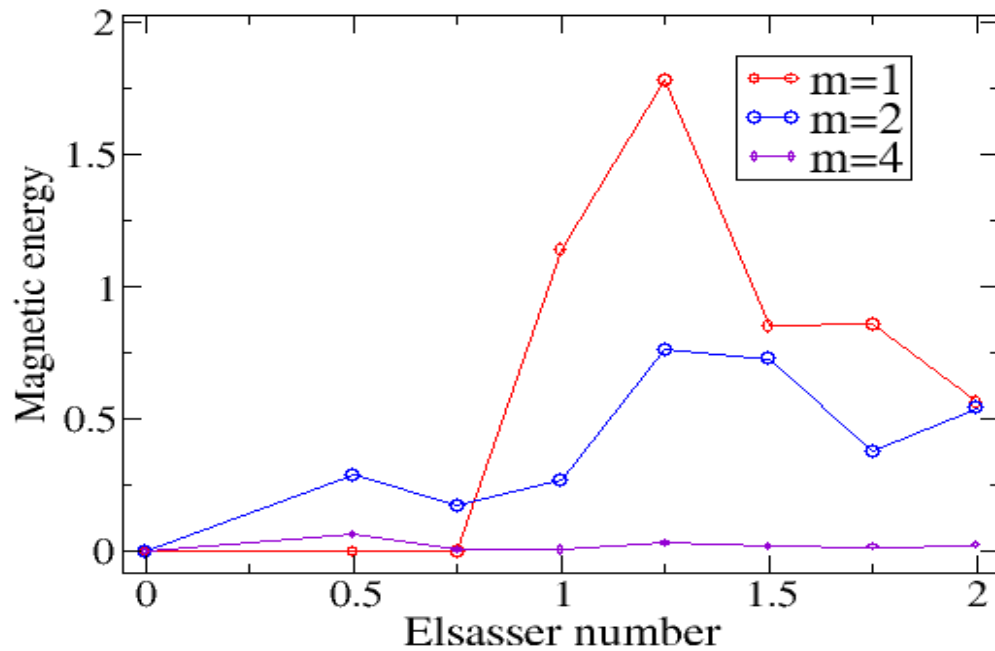
$$\Omega_i = 8 \cdot \Omega_o$$



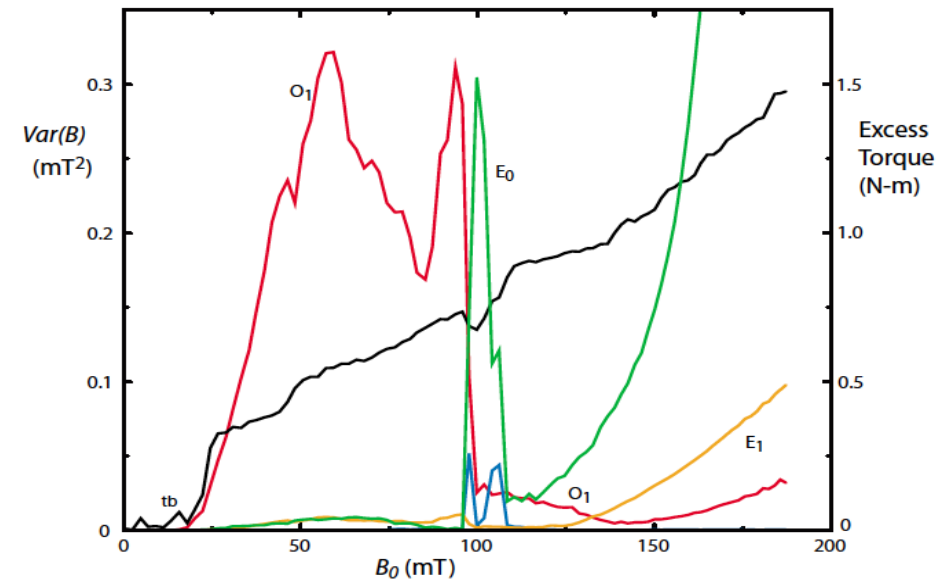
Competition between different modes



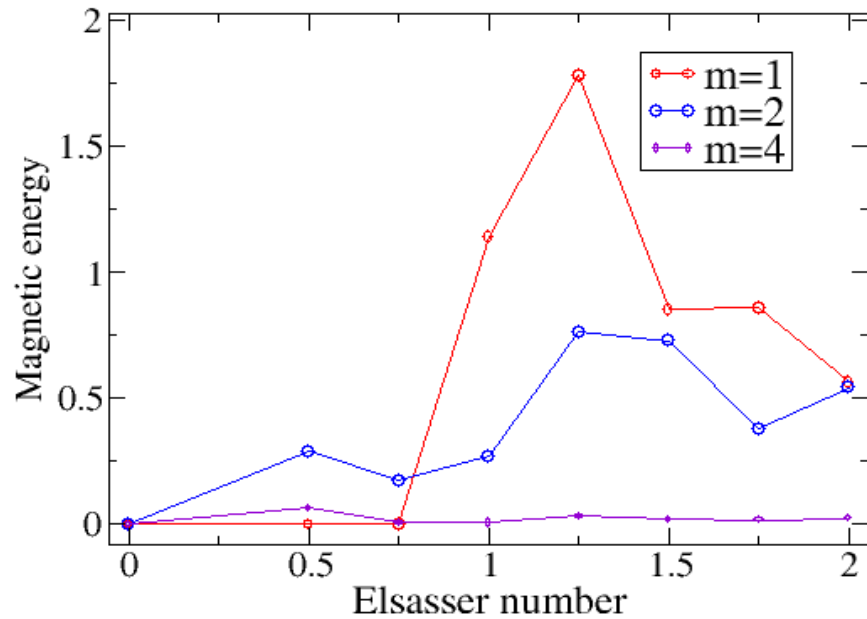
- Transitions between different azimuthal wavenumbers
- Similarities with the Maryland experiment



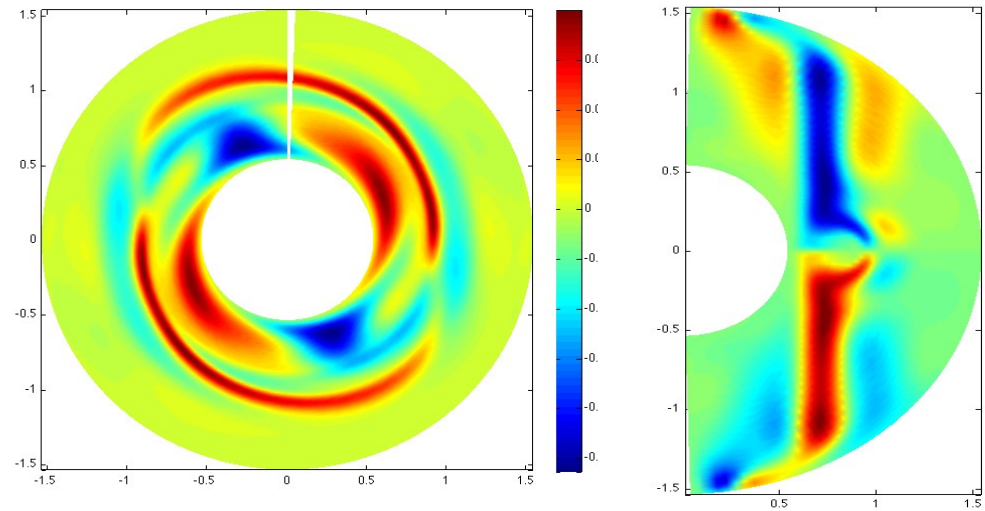
Maryland experiment :



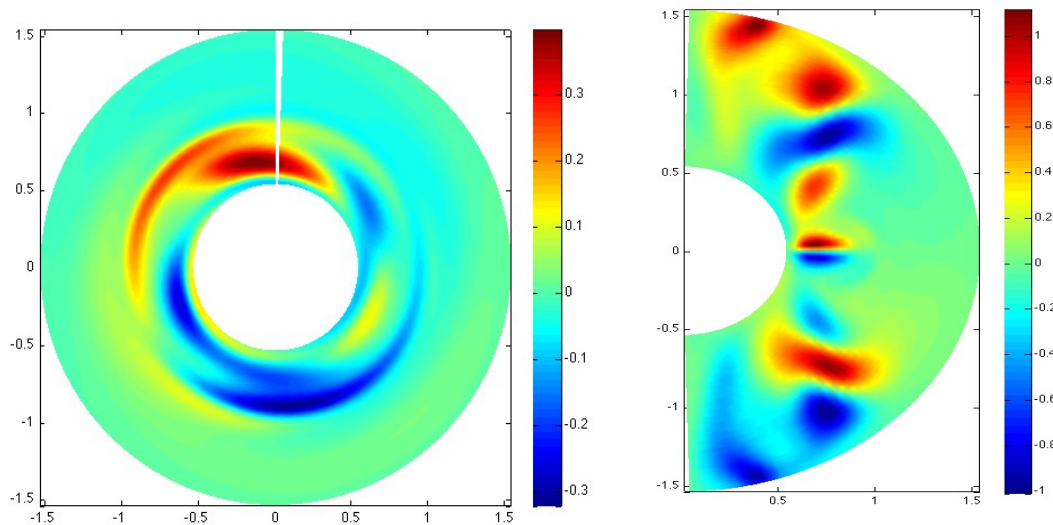
Non-axisymmetric instability



=0.5



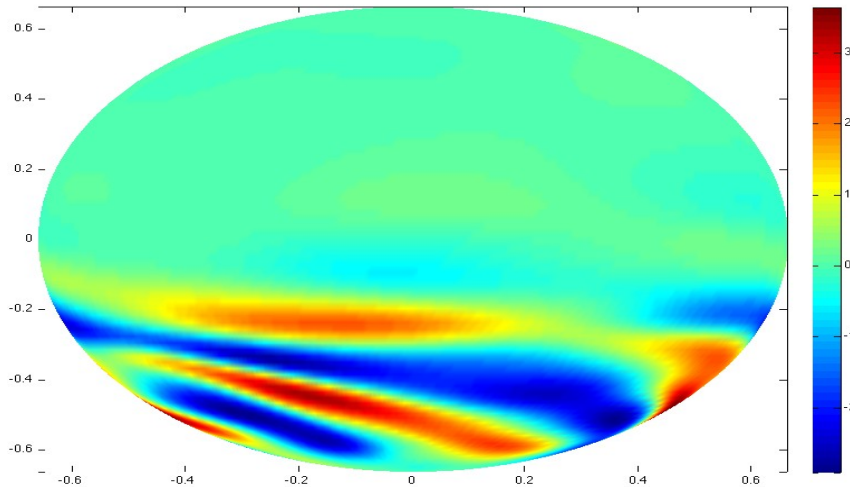
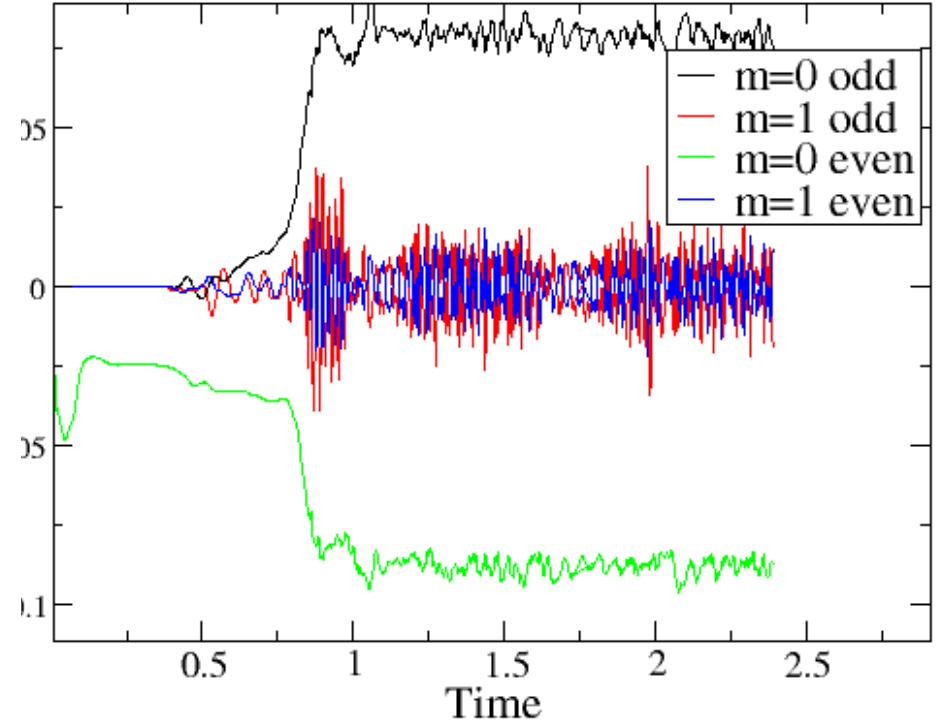
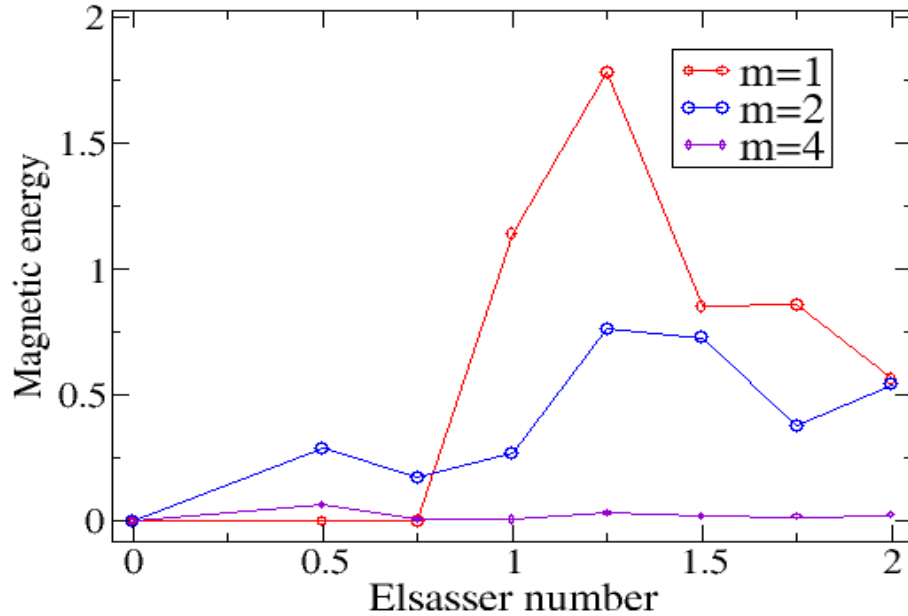
=1.



- ▶ Instability localized close to the tangent cylinder
- ▶ Large vertical wavenumber
- ▶ Interaction symmetric/antisymmetric modes
- ▶ Related to Magnetostrophic MRI (see *Petitdemange 2009*)

Non-linear dynamics

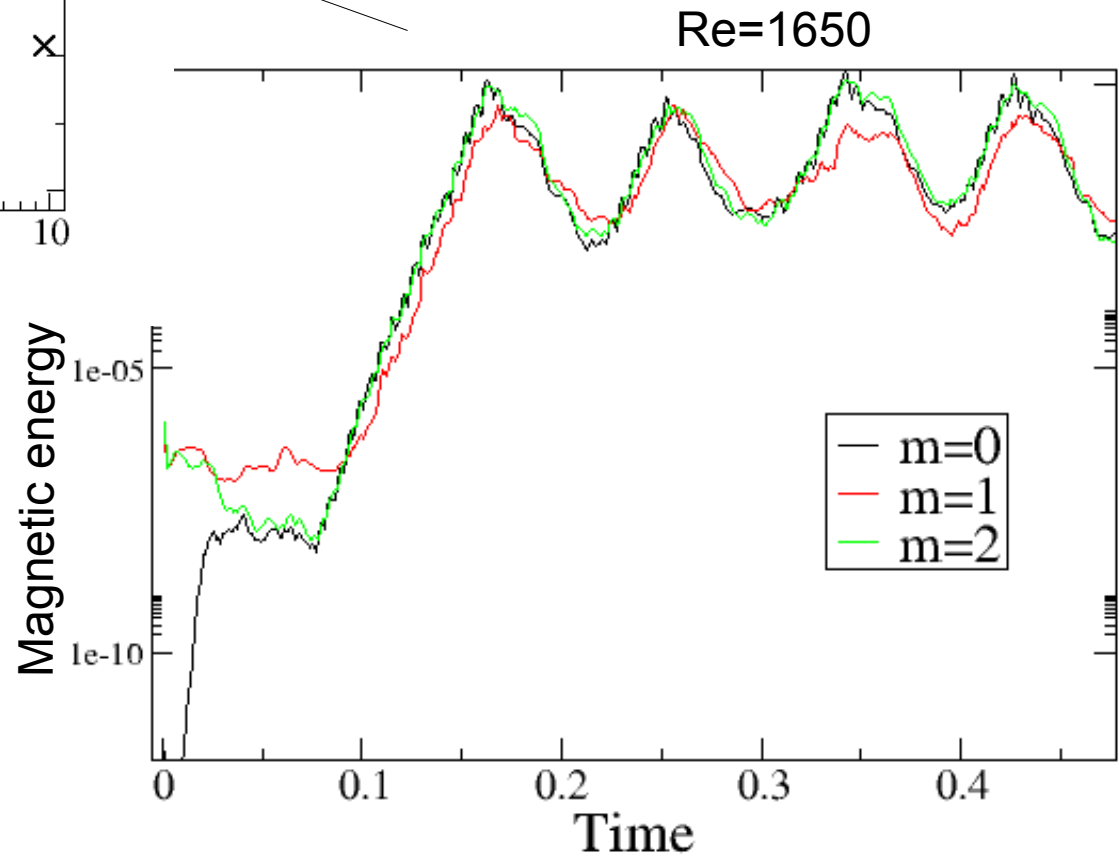
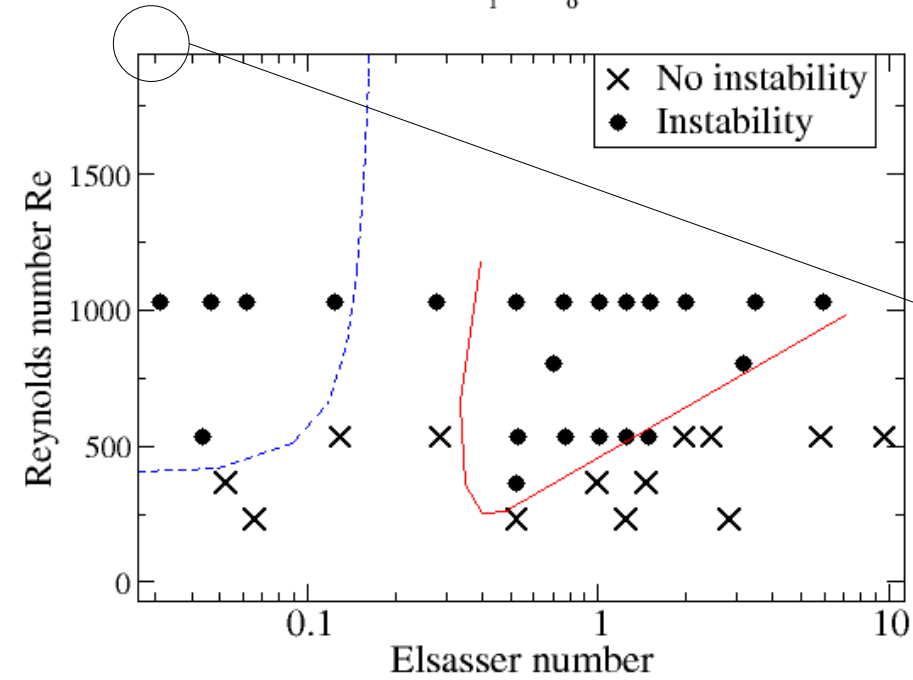
=2



- Hemispherical magnetic field
- Phase locking of symmetric and antisymmetric modes

Interaction with dynamos

$$\Omega_1 = 8 \cdot \Omega_0$$



Conclusion

- Magnetized spherical Couette flow generates non-axisymmetric instabilities
- Non-linear competition between modes
- Similarities with Maryland 'MRI' experiment.

Perspectives

High Reynolds number simulations and effect of Pm

Investigate non-linear dynamics

Comparison with cylindrical simulations

Interaction between MRI and dynamo action