

# Numerical simulations of imbalanced MHD turbulence

Jean C. Perez and Stanislav Boldyrev

University of Wisconsin-Madison  
Department of Physics

April 7, 2009

# Motivation

- Magnetohydrodynamics (MHD) is invoked to describe large scale plasma turbulence in the solar wind, accretion disks, the interstellar medium, etc.

# Motivation

- Magnetohydrodynamics (MHD) is invoked to describe large scale plasma turbulence in the solar wind, accretion disks, the interstellar medium, etc.
- Iroshnikov and Kraichnann 1964 (IK64): Turbulence energy transfer results from nonlinear interaction between counter-propagating Alfvén packets.

# Motivation

- Magnetohydrodynamics (MHD) is invoked to describe large scale plasma turbulence in the solar wind, accretion disks, the interstellar medium, etc.
- Iroshnikov and Kraichnann 1964 (IK64): Turbulence energy transfer results from nonlinear interaction between counter-propagating Alfvén packets.
- The excess of energy in waves moving in one direction is quantified by the cross helicity, which is an ideal invariant.
- When the average cross-helicity vanishes the turbulence is called **Balanced**, otherwise **Imbalanced**.

# Motivation

- Magnetohydrodynamics (MHD) is invoked to describe large scale plasma turbulence in the solar wind, accretion disks, the interstellar medium, etc.
- Iroshnikov and Kraichnann 1964 (IK64): Turbulence energy transfer results from nonlinear interaction between counter-propagating Alfvén packets.
- The excess of energy in waves moving in one direction is quantified by the cross helicity, which is an ideal invariant.
- When the average cross-helicity vanishes the turbulence is called **Balanced**, otherwise **Imbalanced**.
- Recent numerical simulations and observations support the emerging picture that turbulence is locally imbalanced even when the turbulence is globally balanced.

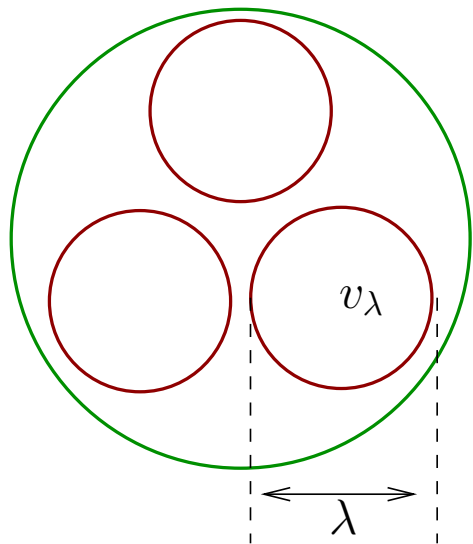
# Motivation

- Magnetohydrodynamics (MHD) is invoked to describe large scale plasma turbulence in the solar wind, accretion disks, the interstellar medium, etc.
- Iroshnikov and Kraichnann 1964 (IK64): Turbulence energy transfer results from nonlinear interaction between counter-propagating Alfvén packets.
- The excess of energy in waves moving in one direction is quantified by the cross helicity, which is an ideal invariant.
- When the average cross-helicity vanishes the turbulence is called **Balanced**, otherwise **Imbalanced**.
- Recent numerical simulations and observations support the emerging picture that turbulence is locally imbalanced even when the turbulence is globally balanced.
- Current understanding of MHD turbulence belongs to the case balanced turbulence. **Turbulence in the Solar Wind is imbalanced!**

# HD vs MHD turbulence pictures

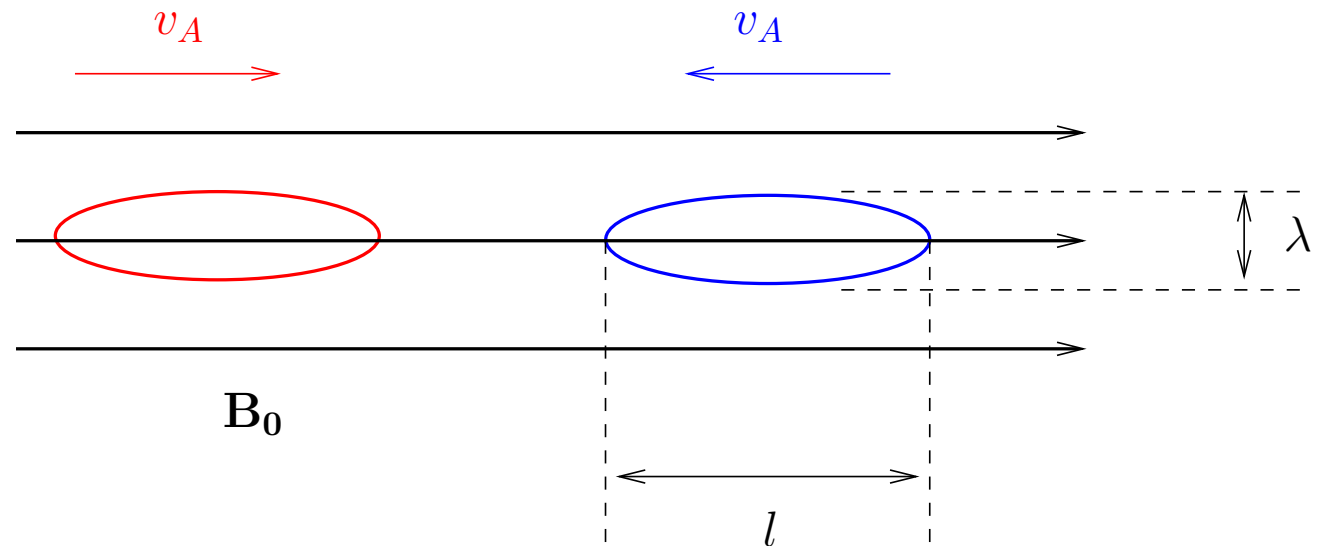
Hydrodynamic turbulence:

- Interaction of eddies.
- Isotropic.
- Strong interactions.
- Energy Cascade.



Magnetohydrodynamics:

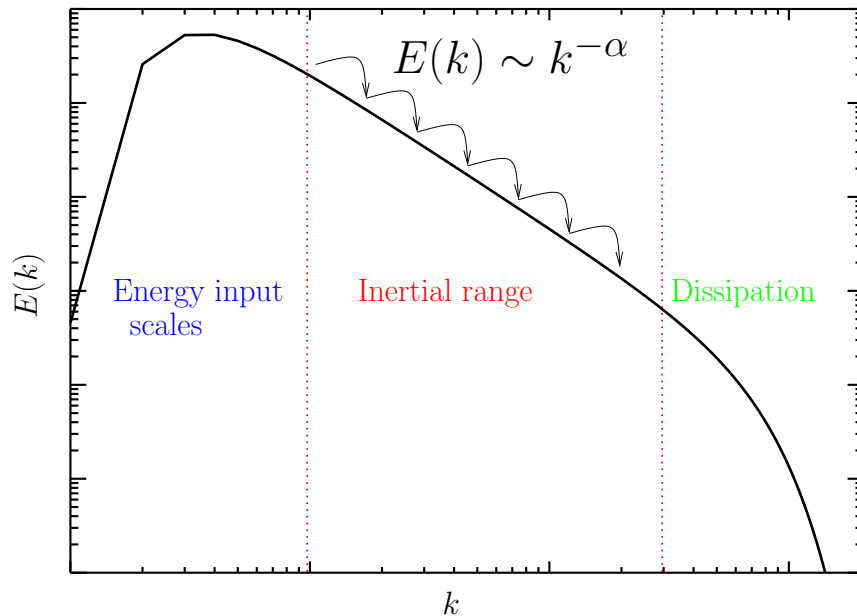
- Interaction of wave packets.
- Anisotropic.
- Weak or Strong interactions.
- Energy and Cross-Helicity cascades.
- Balanced or imbalanced.



# Homogeneous and Isotropic Hydrodynamic (HD) Turbulence

Navier-Stokes equation

$$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} + \nabla p = \frac{1}{R} \nabla^2 \mathbf{v} + \mathbf{f}$$



- For an eddy of size  $\lambda \sim 1/k$  with velocity  $v_\lambda$ , the nonlinear time scale is dimensionally given by

$$\tau_\lambda \sim \frac{\lambda}{v_\lambda}$$

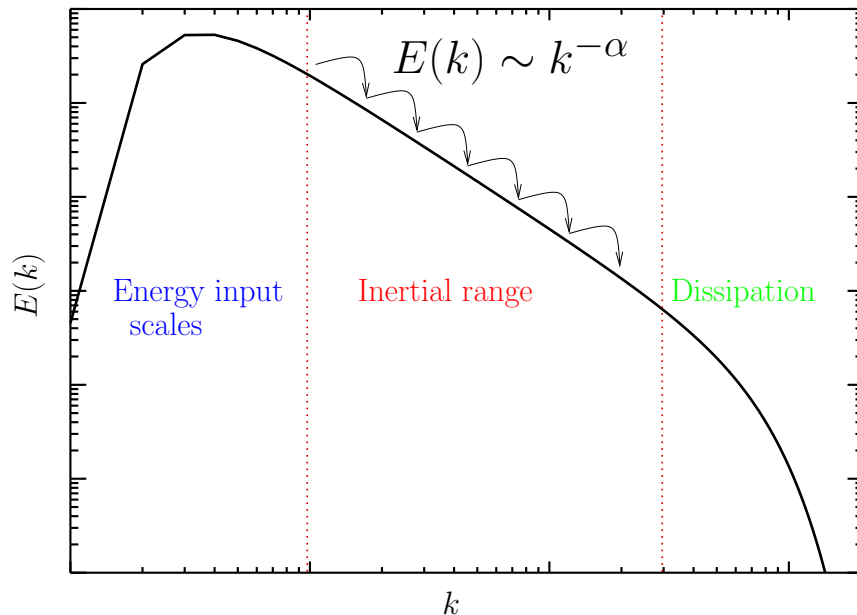
- In steady state

$$\epsilon \sim \frac{v_\lambda^2}{\tau_\lambda} = \text{const} \Rightarrow v_\lambda \sim (\epsilon \lambda)^{1/3}$$

# Homogeneous and Isotropic Hydrodynamic (HD) Turbulence

Navier-Stokes equation

$$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} + \nabla p = \frac{1}{R} \nabla^2 \mathbf{v} + \mathbf{f}$$



- For an eddy of size  $\lambda \sim 1/k$  with velocity  $v_\lambda$ , the nonlinear time scale is dimensionally given by

$$\tau_\lambda \sim \frac{\lambda}{v_\lambda}$$

- In steady state

$$\epsilon \sim \frac{v_\lambda^2}{\tau_\lambda} = \text{const} \Rightarrow v_\lambda \sim (\epsilon \lambda)^{1/3}$$

- Kolmogorov's 5/3 law (K41)

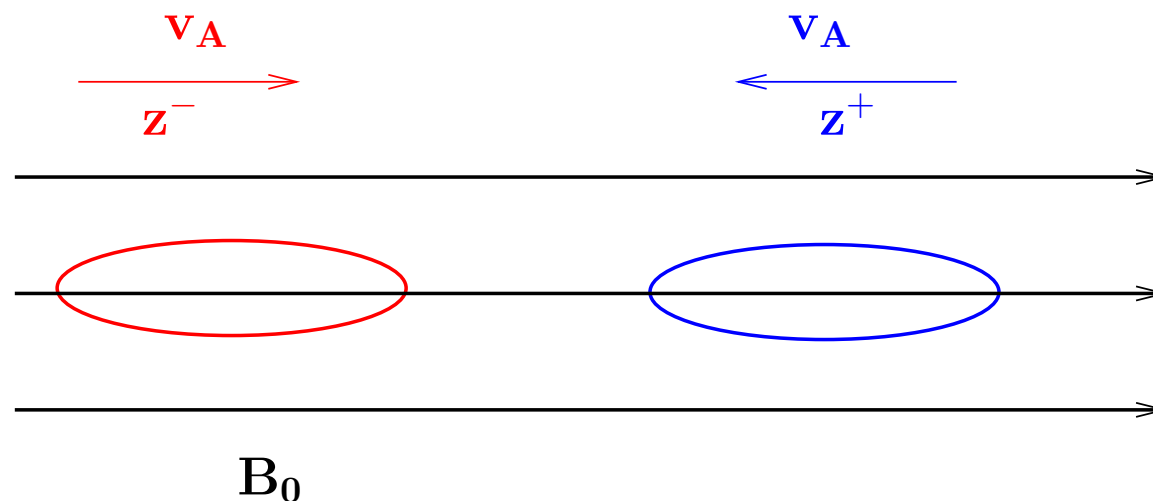
$$E(k) \sim k^{-5/3}, \quad \tau_\lambda \sim k^{-2/3}$$

# Incompressible Magnetohydrodynamic (MHD) turbulence

- MHD turbulence is best studied in the Elsässer formulation

$$\frac{\partial \mathbf{z}^\pm}{\partial t} \mp \underbrace{(\mathbf{v}_A \cdot \nabla) \mathbf{z}^\pm}_{\sim 1/\tau_A} + \underbrace{(\mathbf{z}^\mp \cdot \nabla) \mathbf{z}^\pm}_{\sim 1/\tau_{NL}} = -\nabla P + \frac{1}{R} \nabla^2 \mathbf{z}^\pm + \mathbf{f}^\pm$$

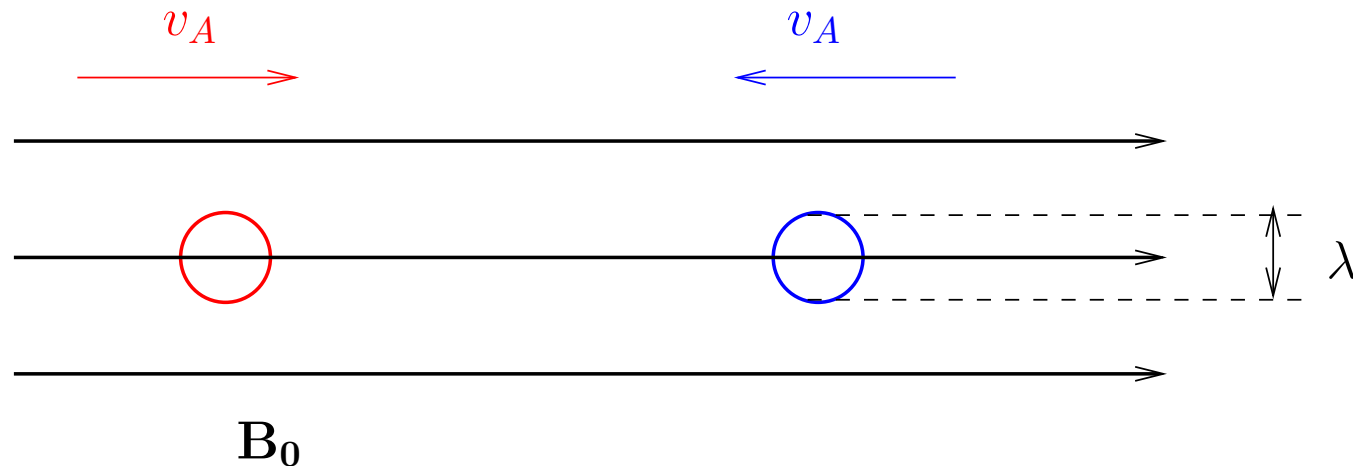
where  $\mathbf{z}^\pm = \mathbf{v} \pm \mathbf{b}$ , and  $\nabla \cdot \mathbf{z}^\pm = 0$ .



- Two competing time scales in the problem: wave time scale  $\tau_A$  and interaction time scale  $\tau_{NL}$ .

# Phenomenological models of Strong Balanced MHD turbulence

- Iroshnikov and Kraichnann (IK): Turbulence is isotropic, energy transfer results from nonlinear interaction between counter-propagating Alfvén packets:

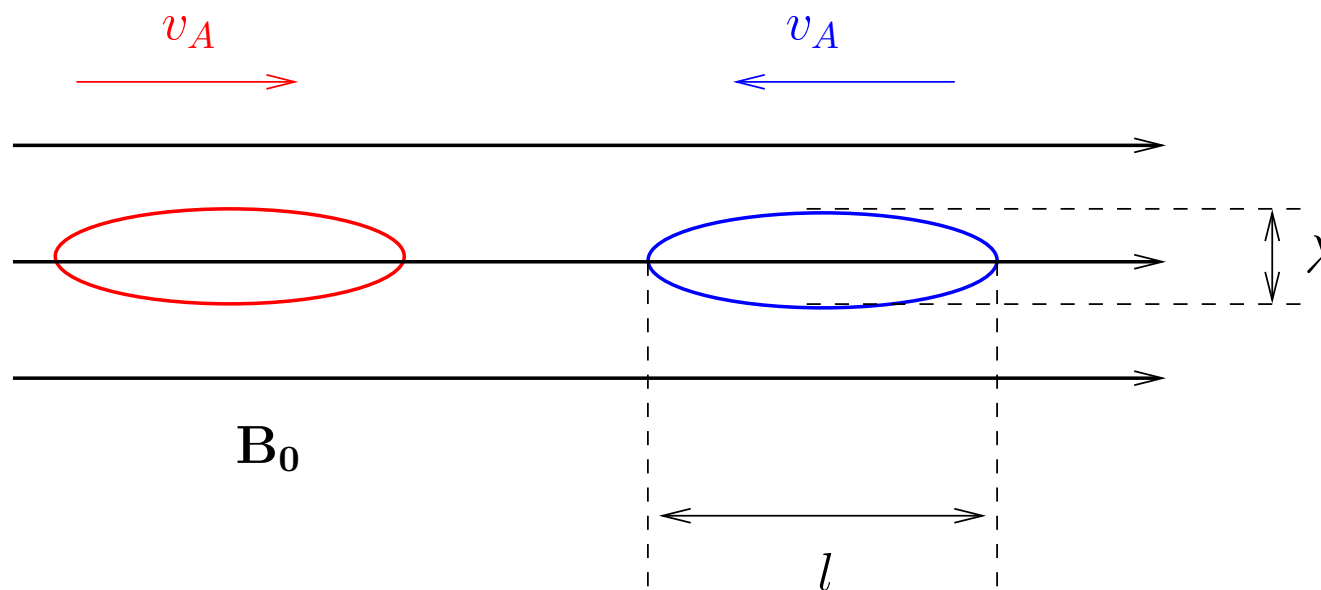


- Energy spectrum:

$$E(k) \sim k^{-3/2}$$

# Phenomenological models of Strong Balanced MHD turbulence

- Goldreich & Sridhar (GS): Eddies become elongated along the local field until there is a formal balance between the linear and nonlinear time scales ( $\tau_A \sim \tau_{NL}$ ), called *Critical Balance*

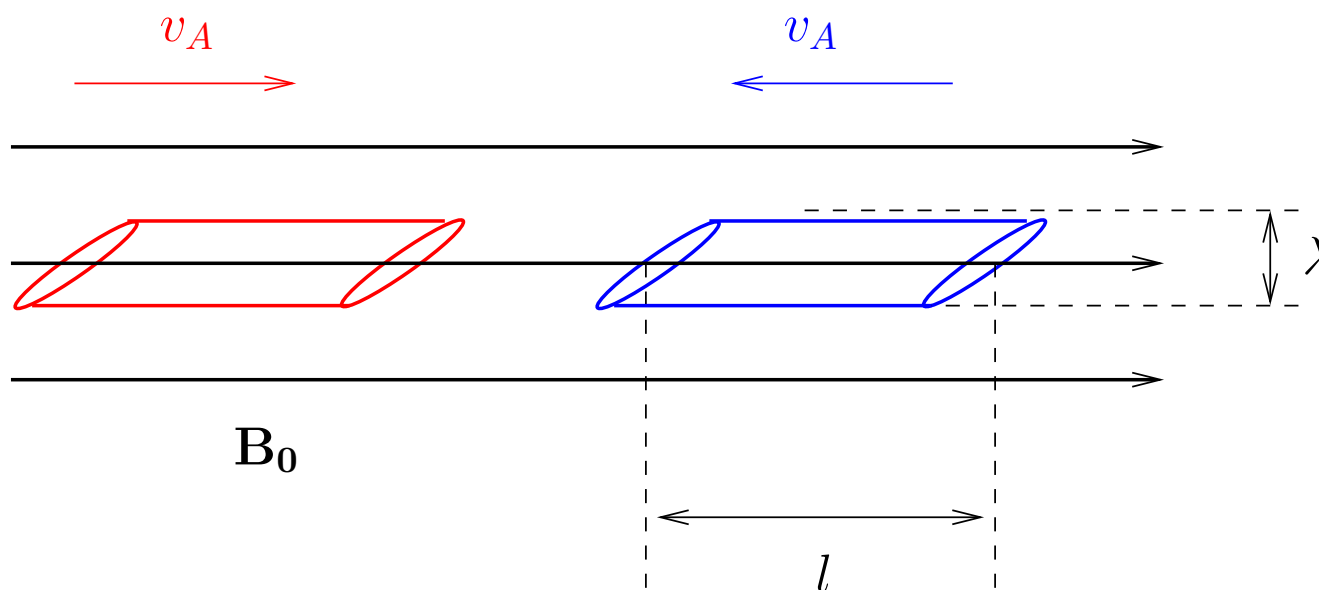


- Energy spectrum:

$$\underbrace{\frac{v_A}{l} \sim \frac{v_\lambda}{\lambda}}_{\text{Critical Balance}} \Rightarrow E(k_\perp) \sim k_\perp^{-5/3}$$

# Phenomenological models of Strong Balanced MHD turbulence

- S. Boldyrev (SB), assumed the nonlinear interaction is depleted due to increasing alignment between velocity and magnetic field fluctuations, while still maintaining critical balance



- Energy spectrum:

$$\frac{v_A}{\lambda_{\parallel}} \sim \frac{v_{\lambda}}{\lambda_{\perp}} \underbrace{\frac{v_{\lambda}}{v_A}}_{\theta_{\lambda}} \Rightarrow E(k_{\perp}) \sim k_{\perp}^{-3/2}$$

# Energy and Cross-Helicity cascades

- Energy and Cross-helicity are defined as

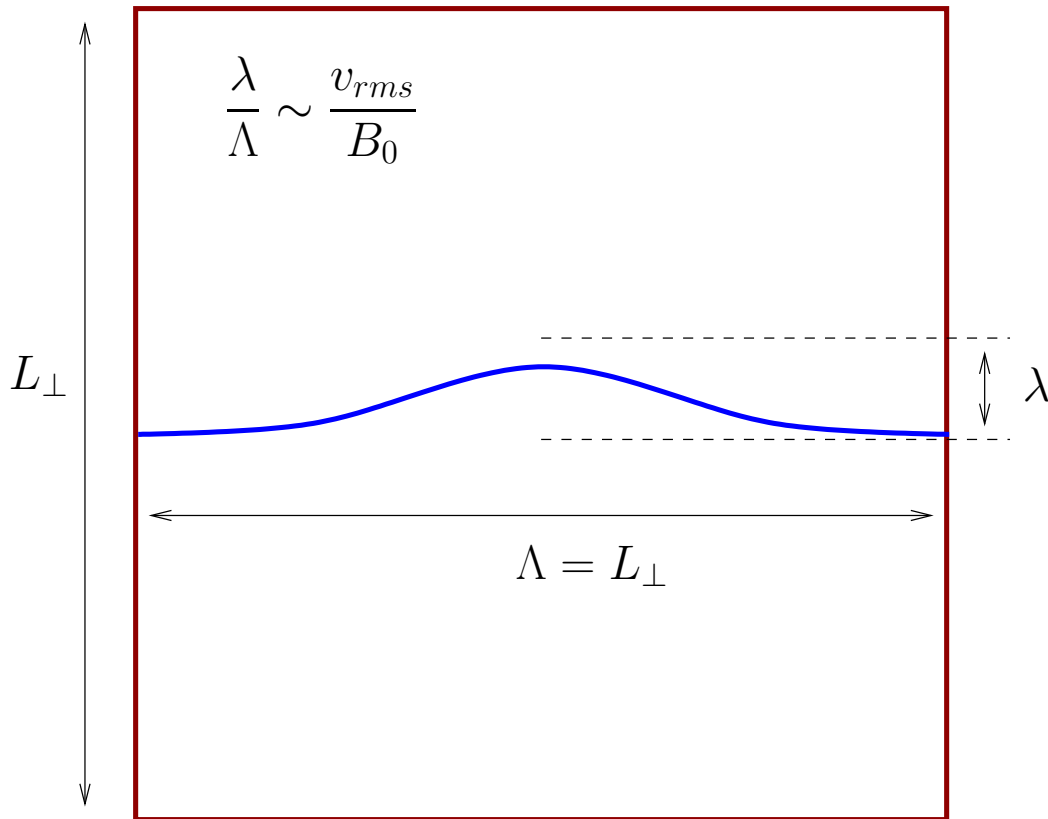
$$E = \frac{1}{2} \langle \mathbf{v}^2 \rangle + \frac{1}{2} \langle \mathbf{b}^2 \rangle, \quad H_c = \langle \mathbf{v} \cdot \mathbf{b} \rangle$$

- $H_c$  measures the degree of correlation between velocity and magnetic fluctuations in the turbulent state.
- In terms of Elsasser variables energy and cross-helicity invariants take the form

$$E = E^+ + E^-, \quad H_c = E^+ - E^-, \quad E^\pm \equiv \frac{1}{4} \langle |\mathbf{z}^\pm|^2 \rangle$$

- Cross-helicity measures the imbalance between counter-propagating waves.
- Together with energy, cross-helicity undergoes a turbulent cascade from large to small scales. The same holds for  $E^\pm$ .

# Anisotropy and computational domain



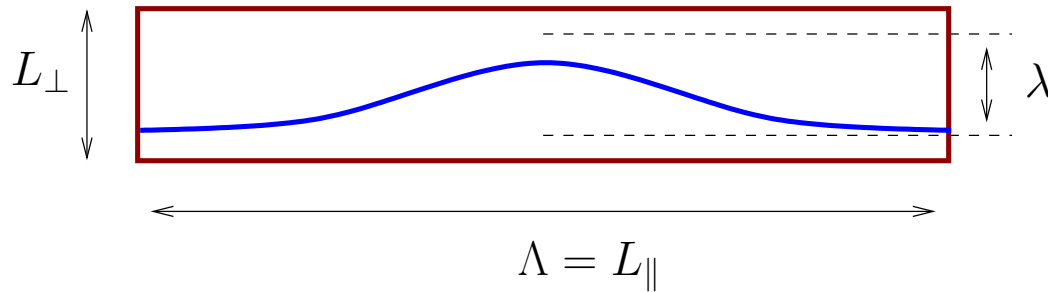
- MHD turbulence is anisotropic: eddies elongated along  $\mathbf{B}_0$ .
- Cubic simulation box is not optimal: a significant number of perpendicular scales wasted.
- Can lead to inaccurate measurements of spectral index.
- The box aspect ratio is crucial to maximize the range of useful perpendicular scales.

- Ratio between linear and nonlinear time scales is controlled at the forcing scale

$$\chi = \frac{k_{\parallel} v_A}{k_{\perp} v_0}$$

# Anisotropy and computational domain

$$\frac{\lambda}{\Lambda} \sim \frac{v_{rms}}{B_0}$$



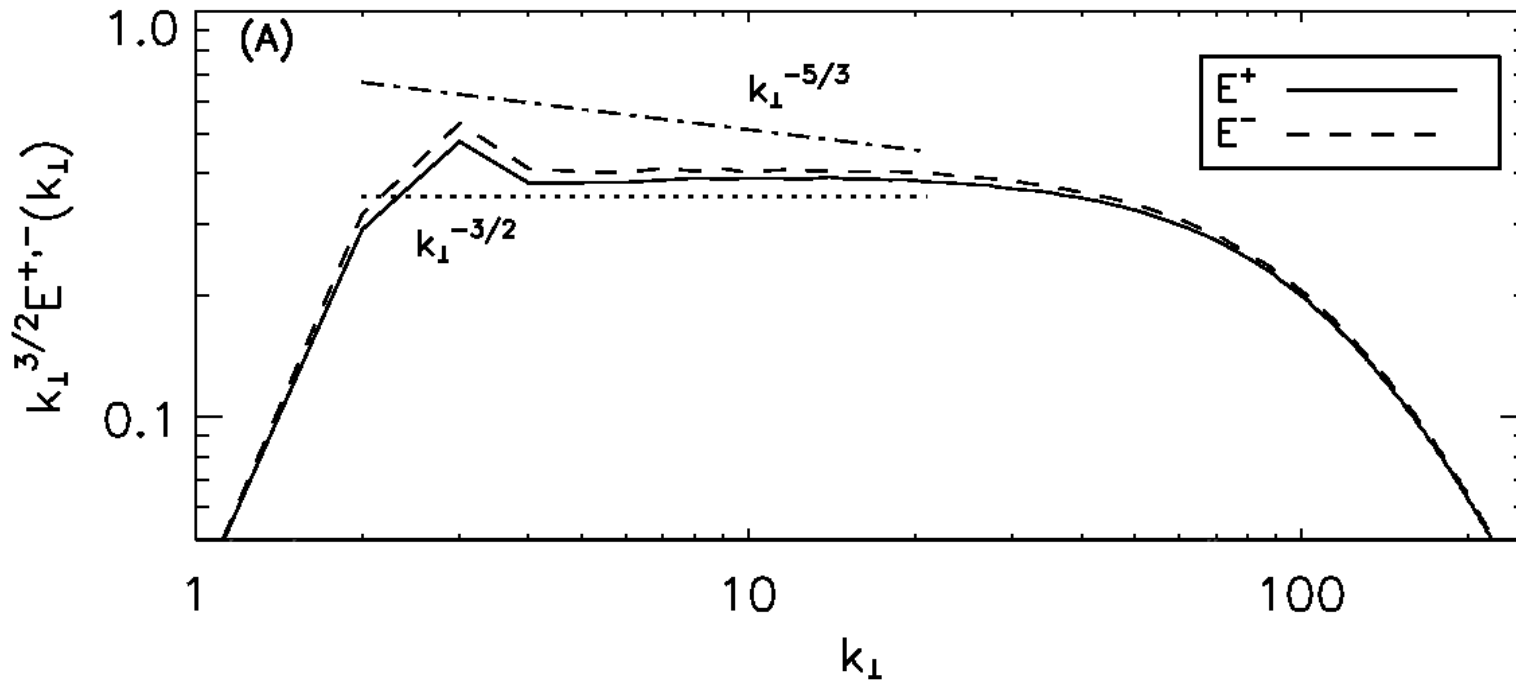
- MHD turbulence is anisotropic: eddies elongated along  $\mathbf{B}_0$ .
  - Cubic simulation box is not optimal: a significant number of perpendicular scales wasted.
  - Can lead to inaccurate measurements of spectral index.
  - The box aspect ratio is crucial to maximize the range of useful perpendicular scales.
- 
- Ratio between linear and nonlinear time scales is controlled at the forcing scale

$$\chi = \frac{k_{\parallel} v_A}{k_{\perp} v_0}$$

## Reduced MHD turbulence code

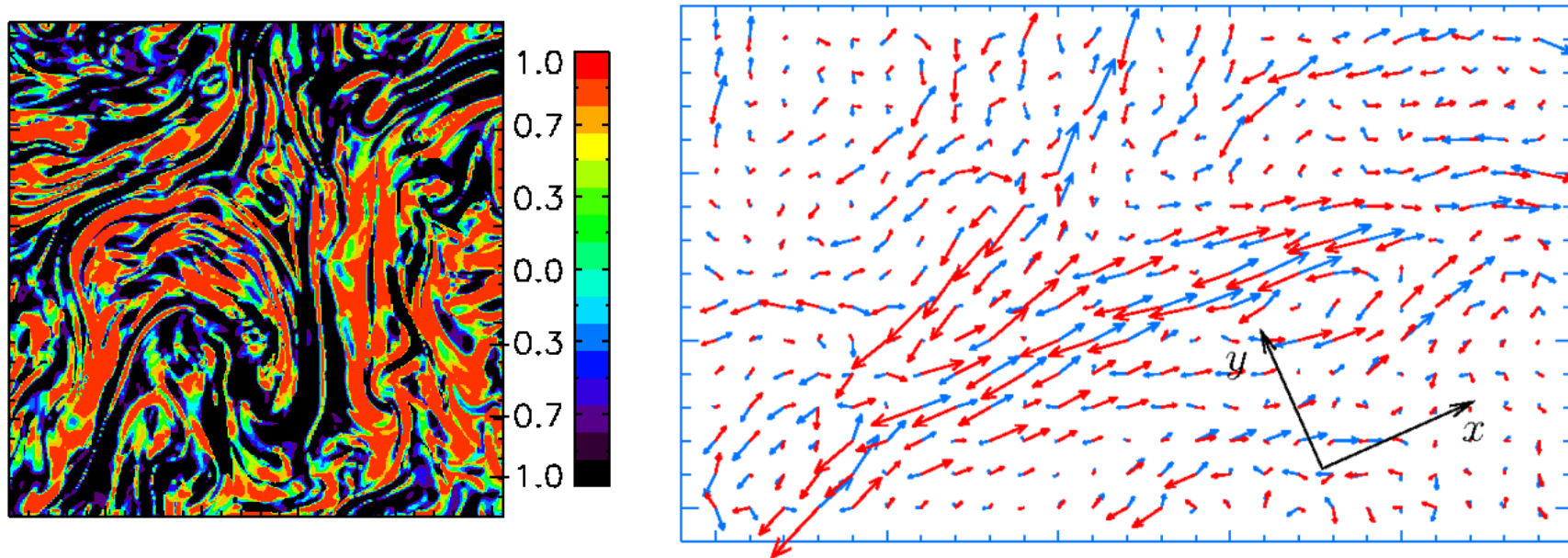
- We used a 2/3-dealiased pseudo-spectral method to solve RMHD equations.
- No artificial viscosity is used (hyper-viscosity), to minimize undesired bottleneck effects.
- Simulations are regularly performed on 1024 to 2048 processors in Ranger supercomputer at the Texas Advanced Computing Center, UT-Austin through an NSF Teragrid Allocation.
- Code has good strong and weak scaling up 16,384 cores on ranger.
- We have used over 2 million CPU hours, it would roughly take more than two centuries on a single computer.

# Balanced Turbulence Simulations on a $1024^2 \times 256$ grid



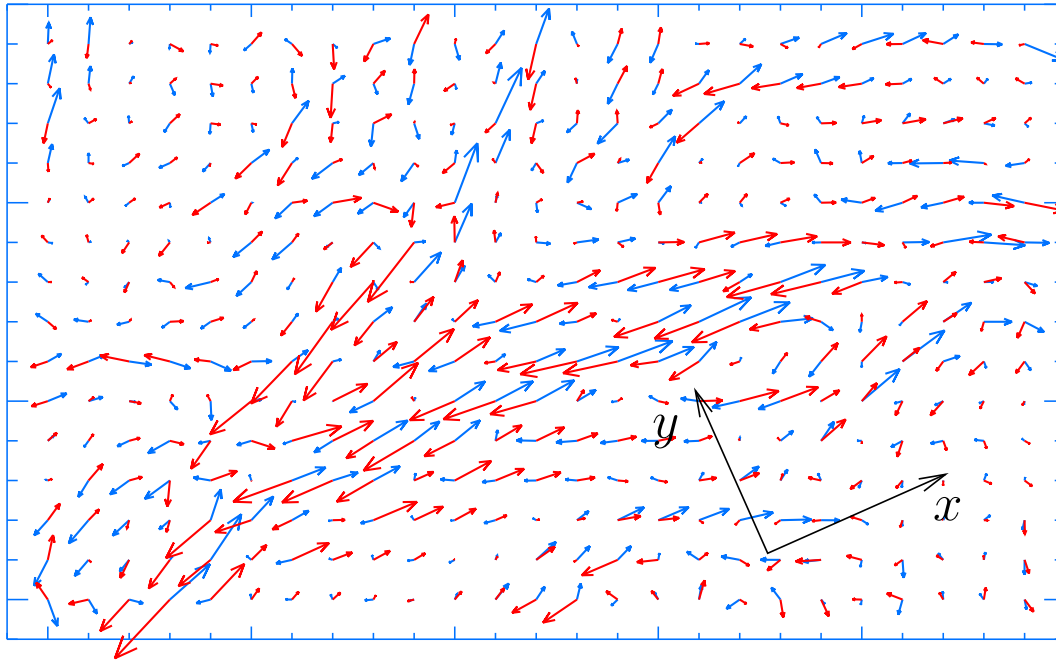
- Results for  $(256^2, 512^2$  and  $1024^2) \times 256$ , are all consistent with  $k_{\perp}^{-3/2}$ . (No bottleneck is observed!)
- Consistent with numerics by: Maron & Goldreich 2001, Haugen, Brandenburg & Doubler 2003, Muller & Grappin 2005, Mason, Cattaneo & Boldyrev 2008, Mininni & Pouquet 2008.

# The role of cross helicity $H_c = E^+ - E^-$



- Left panel shows cosine of the alignment angle between  $\mathbf{v}_\lambda$  and  $\mathbf{b}_\lambda$  fluctuations in the guide-field perpendicular plane at scales  $\lambda = L_\perp/12$  for balanced turbulence. Right panel shows velocity and magnetic fields at one alignment region.
- Domains are split in patches of highly aligned and anti-aligned regions.
- A theory of imbalanced turbulence should be consistent with the fact that **balanced turbulence** is composed of imbalanced regions.

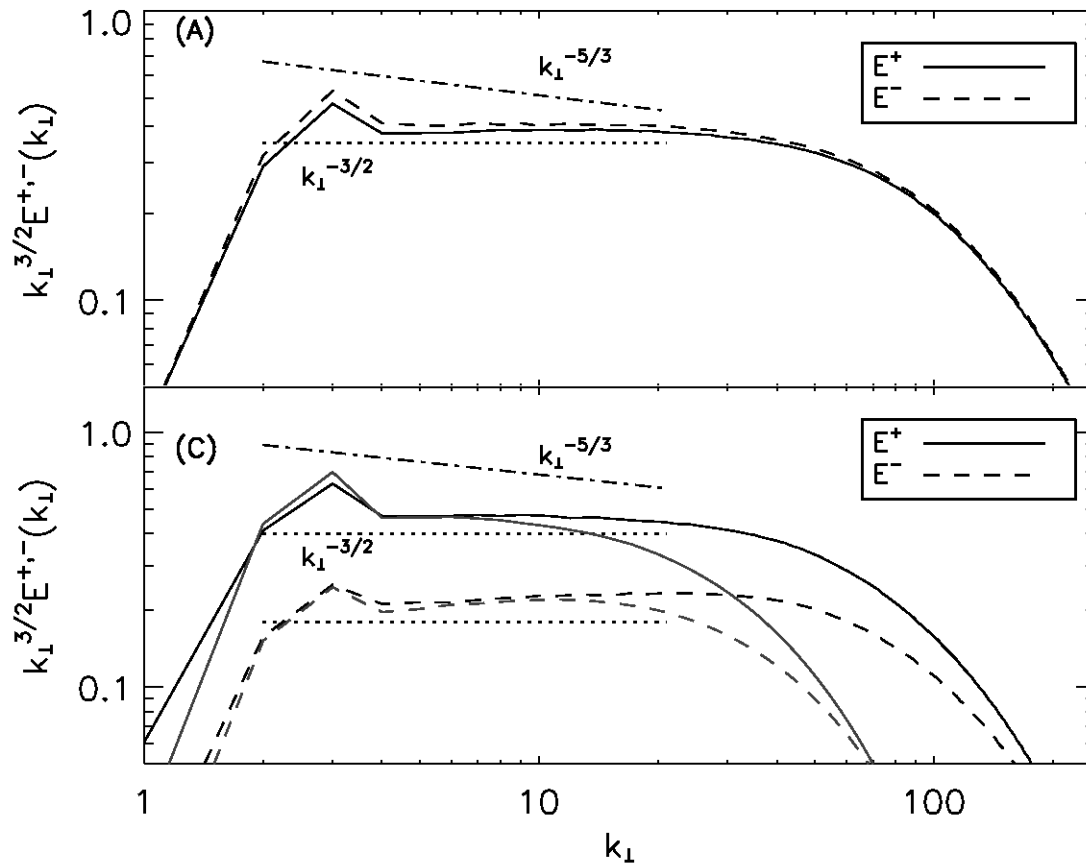
# Dynamic alignment model for the observed spectra: balanced case



- Fluctuations of velocity and magnetic field develop correlated regions of alignment and anti-alignment.
- Fluctuations are aligned toward the  $x$  direction, while gradient of fluctuations is stronger in the perpendicular direction ( $y$ ).
- Alignment is scale dependent. When the angle decreases with the scale as  $\theta_\lambda \sim \lambda^{1/4}$  the constant energy flux constraint leads to the energy scaling  $E(\mathbf{k}_\perp) \sim k_\perp^{-3/2}$ .

- Lithwick and Goldreich 2007 argue that energy cascade times are different for  $z^+$  and  $z^-$  and lead to scalings  $E^+ \sim E^- \sim k_{\perp}^{-5/3}$ .
- Beresnyak & Lazarian 2008, and Chandran 2008 assume that the high amplitude wave, say  $z^+$  undergoes a weak cascade while  $z^-$  undergoes a strong cascade, therefore energy spectra  $E^+$  and  $E^-$  have different scalings.
- The fact that balanced MHD turbulence is constituted by regions of both positive and negative cross-helicity is incompatible with current models of imbalanced turbulence.

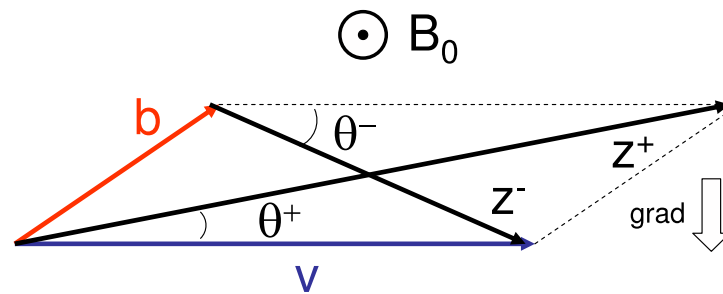
# Imbalanced Turbulence Simulations on a $1024^2 \times 256$ grid



1. Strong Balanced turbulence:  $E^+(\mathbf{k}_{\perp}) \sim E^-(\mathbf{k}_{\perp}) \propto k_{\perp}^{-3/2}$ .
2. Strong Imbalanced turbulence:  $E^+(\mathbf{k}_{\perp}) \propto E^-(\mathbf{k}_{\perp}) \propto k_{\perp}^{-3/2}$ .
3. Independently observed in solar wind data by Podesta & Bhattacharjee 2008, unpublished

# Dynamic alignment model for the observed spectra: imbalanced case

- We assume that alignment is also present in the imbalanced case.
- However, alignment angle is different for  $z^+$  and  $z^-$ , with the geometric constraint  $z^+ \theta^+ \sim z^- \theta^-$ .



- Nonlinear interaction time is the same for both,  $\tau^+ = \tau^-$ .

$$\frac{\partial}{\partial t} \mathbf{z}^{\pm} \mp v_A \partial_{\parallel} \mathbf{z}^{\pm} + \underbrace{(\mathbf{z}^{\mp} \cdot \nabla)}_{z_{\lambda}^{\mp} \theta_{\lambda}^{\mp} k_{\perp} \sim 1/\tau^{\pm}} \mathbf{z}^{\pm} = 0$$

- Constant flux assumption  $(z^{\pm})^2 / \tau_{\lambda}^{\pm} = \text{const}$  leads to same scaling  $k_{\perp}^{-3/2}$ , although  $E^{\pm}$  have different amplitudes.

## Concluding Remarks

- MHD turbulence, both balanced and imbalanced, is characterized by regions of positive and negative cross-helicity.
- Numerical simulations show that the spectra of  $E^+$  and  $E^-$  approach the scaling  $k_{\perp}^{-3/2}$  with different amplitudes.
- Simulations provide a coherent picture of balanced and imbalanced MHD turbulence.